

## Motivation

- Diffusion processes are widespread in many natural processes.
- Typically characterized by three terms: a drift term due to a potential field, the interaction with other particles, and a stochastic term.
- Existing methods either rely on having access to particles trajectories (often not possible) or are data and compute inefficient.

## The JKO scheme

The Fokker-Planck equation,

$$\frac{\partial \rho(t, x)}{\partial t} = \nabla \cdot (\nabla V(x) \rho(t, x)) + \beta \nabla^2 \rho(t, x),$$

describes the time evolution of the distribution  $\rho$  of a set of particles undergoing drift and diffusion,

$$dX(t) = -\nabla V(X(t))dt + \sqrt{2\beta}dW(t),$$

where  $X(t)$  is the state of the particle,  $V(x)$  the driving potential, and  $W(t)$  the Wiener process. It turns out the resulting particles trajectory can be described via the JKO scheme:

$$\mu_{t+1} = \operatorname{argmin}_{\mu \in \mathcal{P}(\mathbb{R}^d)} J(\mu) + \frac{1}{2\tau} W_2(\mu, \mu_t)^2, \quad (1)$$

where  $J$  is an energy functional and  $\tau > 0$  is the time discretization.

### A general energy functional:

$$J(\mu) = \underbrace{\int_{\mathbb{R}^d} V(x) d\mu(x)}_{\text{potential}} + \underbrace{\int_{\mathbb{R}^d \times \mathbb{R}^d} U(x-y) d(\mu \times \mu)(x, y)}_{\text{interaction}} + \beta \underbrace{\int_{\mathbb{R}^d} \rho(x) \log(\rho(x)) dx}_{\text{internal}}$$

## Intuition in $\mathbb{R}^d$

Consider the analog of (1) in the  $\mathbb{R}^d$ ,

$$x_{t+1} = \operatorname{argmin}_{x \in \mathbb{R}^d} J(x) + \frac{1}{2\tau} \|x - x_t\|^2.$$

We replace the above by its first-order optimality condition

$$\nabla J(x_{t+1}) + \frac{1}{\tau}(x_{t+1} - x_t) = 0.$$

Given a dataset  $(x_0, x_1, \dots, x_T)$ , we find  $J$  as:

$$\min_J \sum_{t=0}^{T-1} \left\| \nabla J(x_{t+1}) + \frac{1}{\tau}(x_{t+1} - x_t) \right\|^2.$$

Our work enables an analogous result in the Wasserstein space.

## First-order optimality conditions for the JKO scheme

Our analysis suggest that given a populations dataset  $(\mu_0, \mu_1, \dots, \mu_T)$  we seek the parameters  $\theta$  minimizing

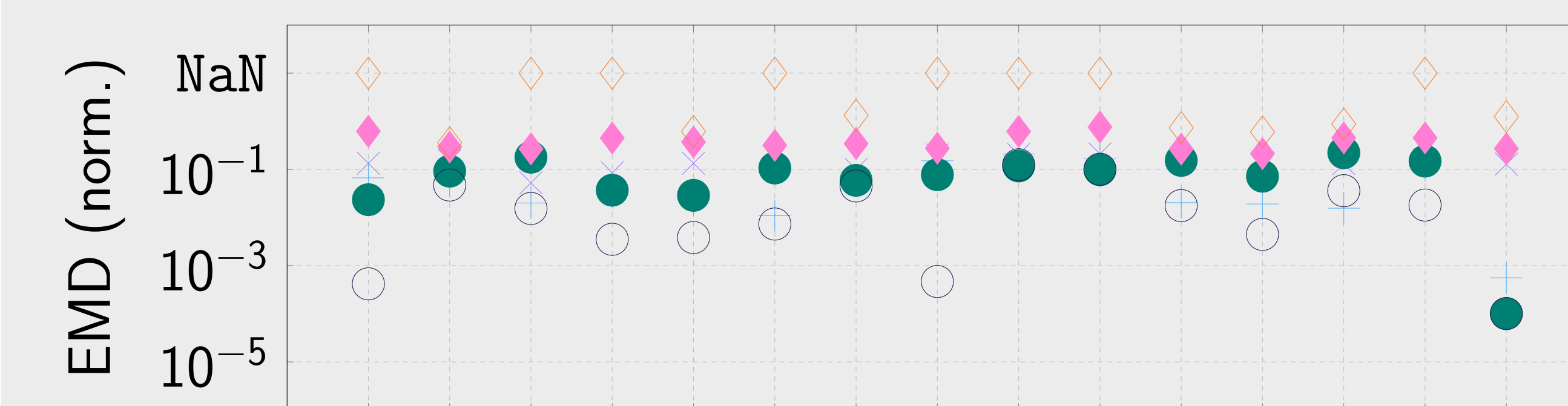
$$\sum_{t=0}^{T-1} \int_{\mathbb{R}^d \times \mathbb{R}^d} \left\| \nabla V_\theta(x_{t+1}) + \int_{\mathbb{R}^d} \nabla U_\theta(x_{t+1} - x'_{t+1}) d\mu_{t+1}(x'_{t+1}) + \beta_\theta \frac{\nabla \rho_{t+1}(x_{t+1})}{\rho_{t+1}(x_{t+1})} + \frac{1}{\tau}(x_{t+1} - x_t) \right\|^2 d\gamma_t(x_t, x_{t+1}), \quad (2)$$

where  $\gamma_0, \gamma_1, \dots, \gamma_T$  are optimal transport plans between  $\mu_t$  and  $\mu_{t+1}$ .

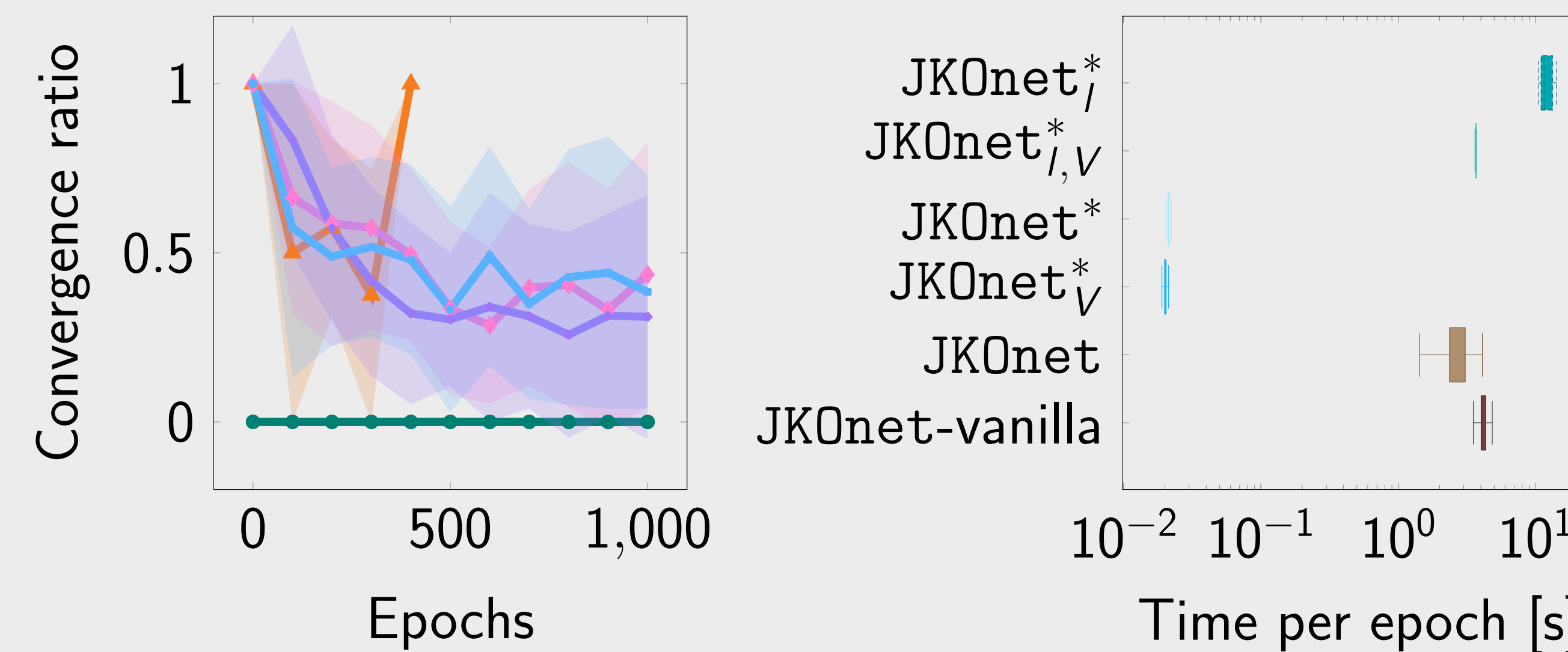
- $\gamma_t$  can be computed once, before-hand, and efficiently;
- when  $J_\theta$  is a neural network, we minimize (2) via gradient descent;
- when  $J_\theta$  is parametrized linearly, we have a closed-form solution.

## Training at lightspeed

JKOnet\* + JKOnet\*\_V ● JKOnet\*\_I ○ JKOnet\*\_{I,V} ◆ JKOnet ◇ JKOnet-vanilla

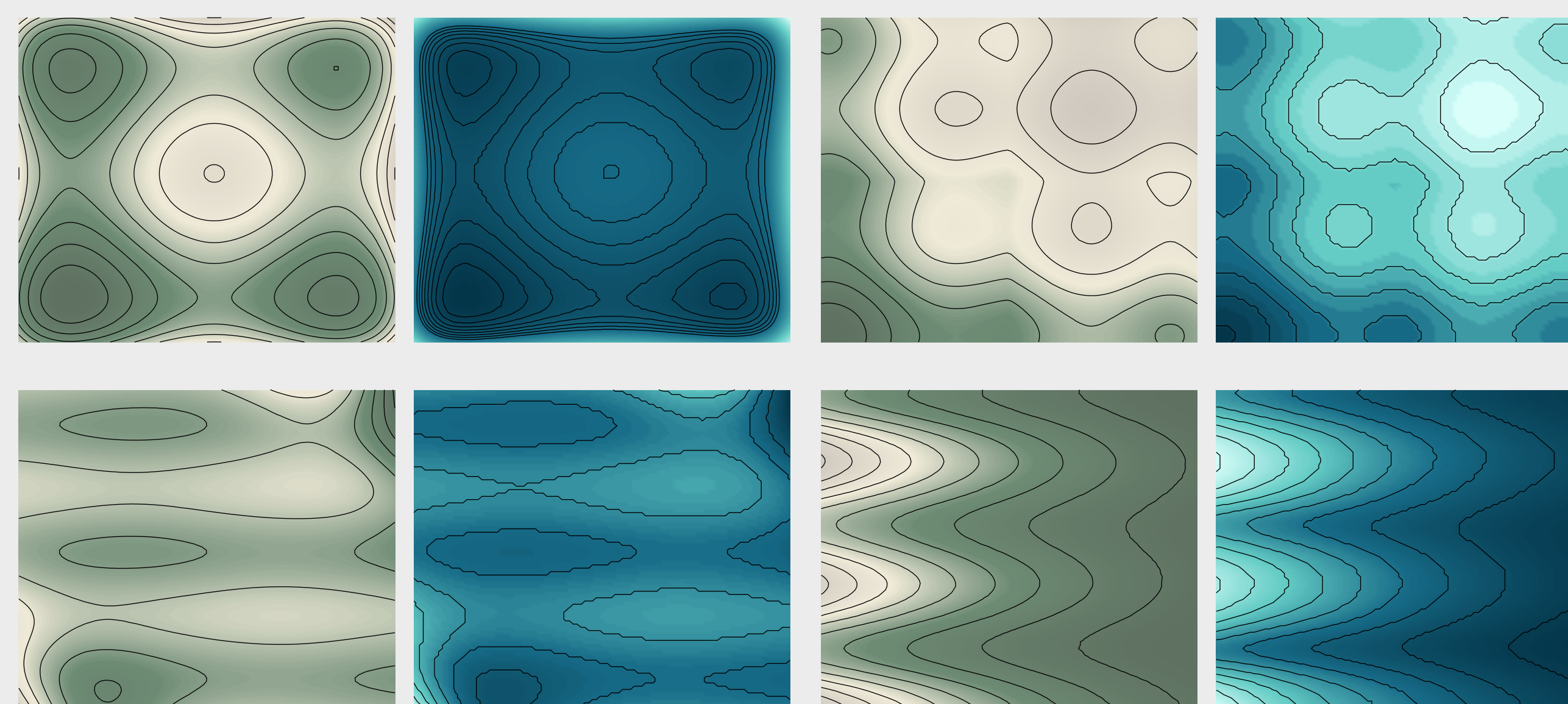


Different potential  $V(x)$  driving the diffusion.



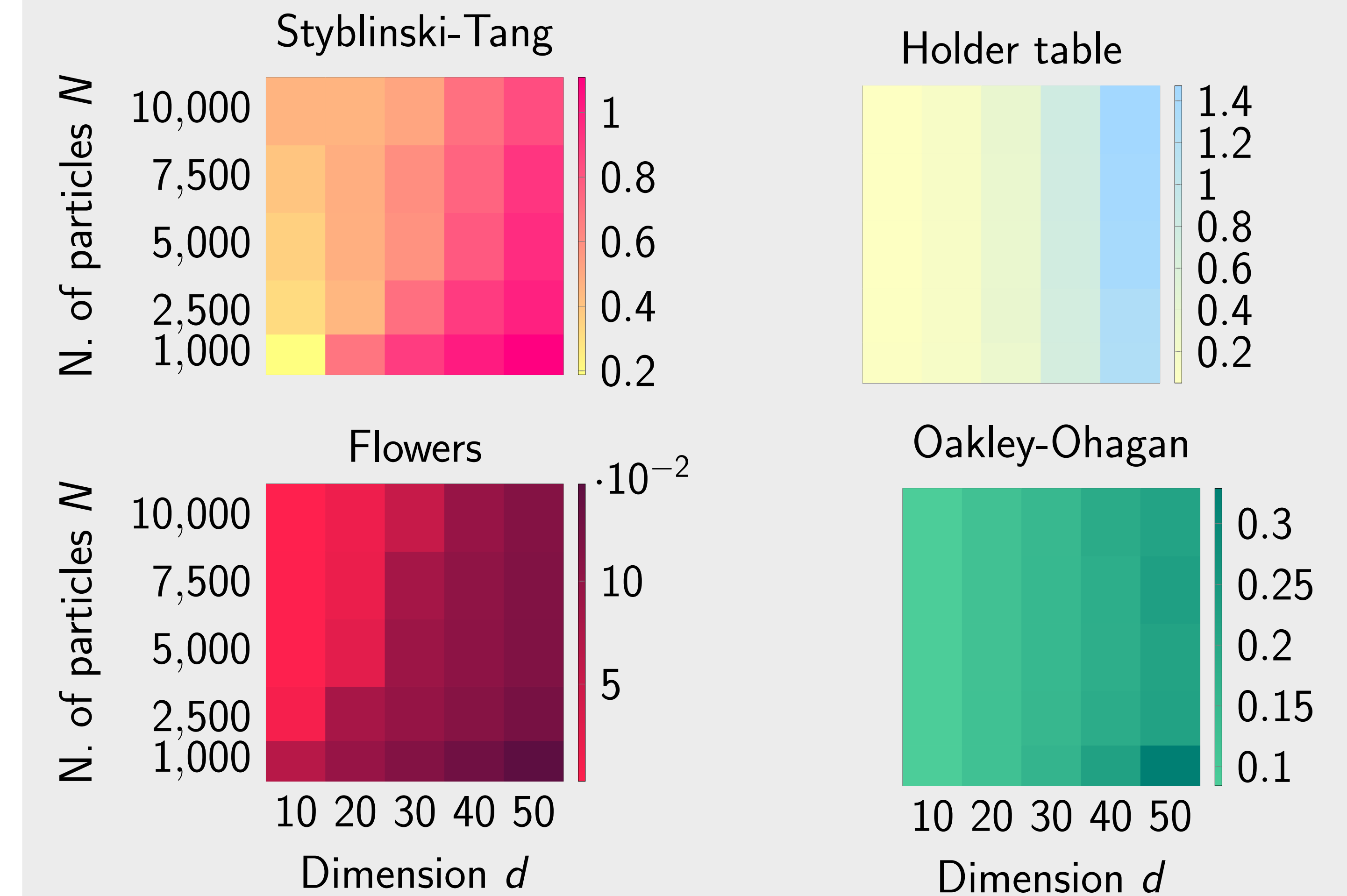
## Eye-candies

Level curves of different potentials (green: real; blue: estimated).



## Scaling to high-dimensions

- Sub-linear error growth.
- Negligible increase in training time.



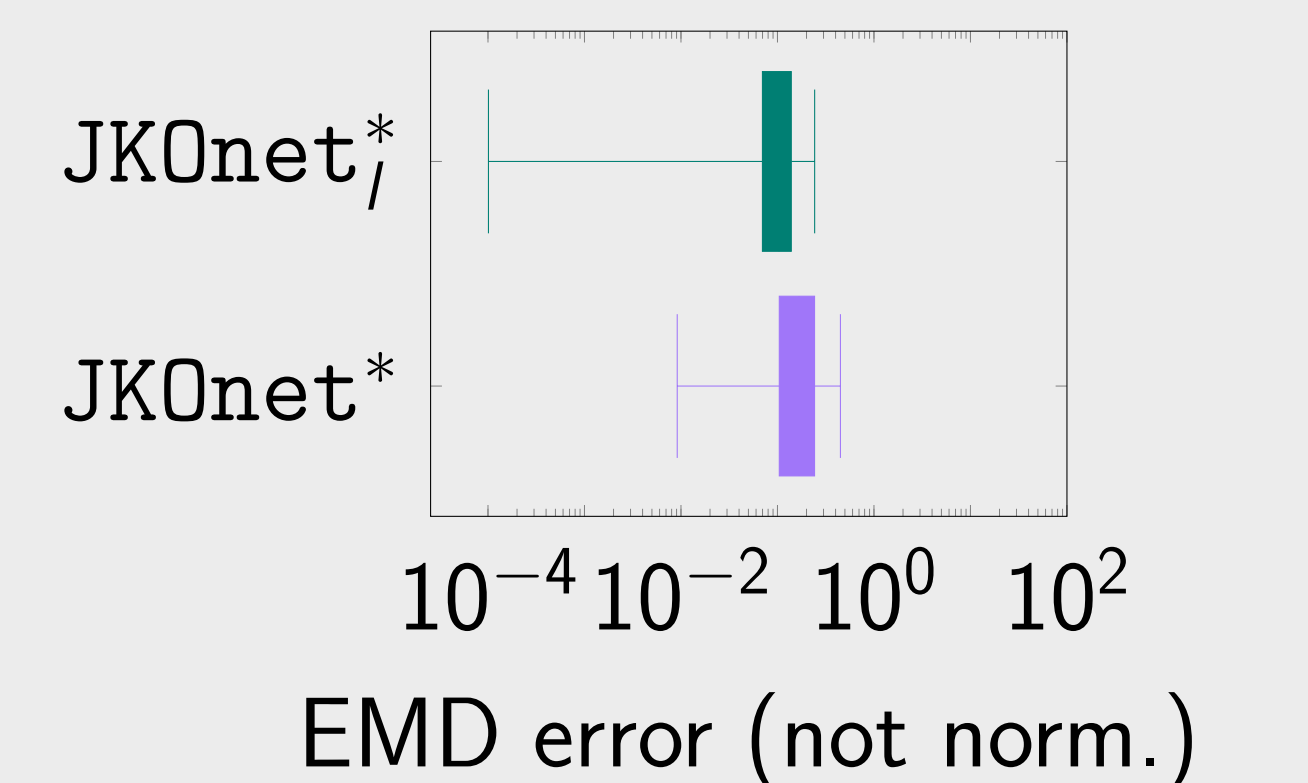
## Learning general energy terms

Additional error sources:

- Sampling error (internal energy)
- Estimation of the densities

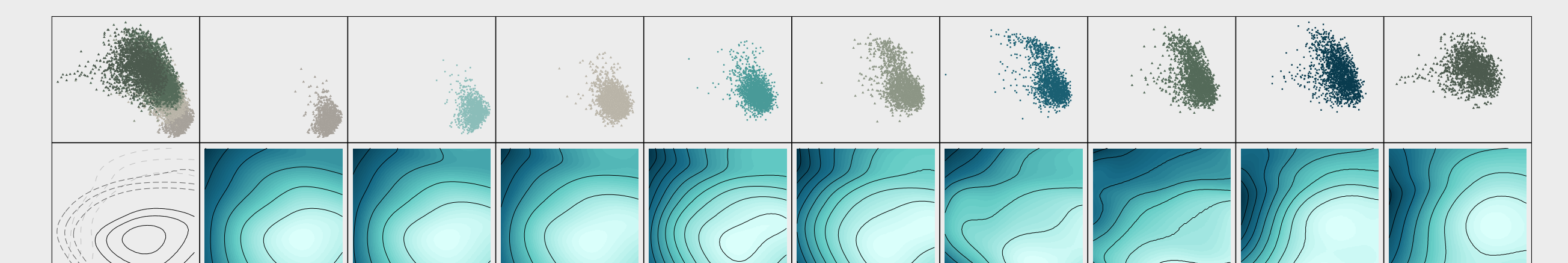
Nonetheless: our method (JKOnet\*, non-linear and linearly parametrized) recovers all the energy terms.

Open question: observability of the different energy terms?



## Learning single-cell diffusion dynamics

- Intuition: cellular evolution minimizes (some) energy;
- We can account for unobserved variables via time-varying energies;
- State-of-the-art accuracy at a fraction of the computational cost of the other methods: less than minute vs hours.



## Cool stuff worth looking at next

- Fast distillation of diffusion models for one- or few-steps generation;
- Observability of the different energy terms;
- Can optimality conditions in  $\mathcal{P}(\mathbb{R}^d)$  be helpful in your work? Reach out!

